

# Use of Carleman estimates for stability in some inverse problems

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**CIRM - Février 2009**

Colloque "Contrôle et Problèmes Inverses pour les EDP : Aspects Théoriques et Numériques"

## ▶ Introduction

- ▶ Usual Schrödinger and wave equations
- ▶ Non-linear inverse problem
- ▶ Uniqueness and Stability
- ▶ Transmission equations

## ▶ Method to prove stability for some linear inverse problem

## ▶ Carleman estimates

- ▶ Weight function
- ▶ Proof scheme of the usual case
- ▶ Transmission case

## ▶ Stability Theorem

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# Schrödinger and wave equations

$T > 0$ ,  $\Omega \subset \mathbb{R}^2$  bounded domain

Smooth boundary  $\partial\Omega$

potential  $q = q(x) \in L^\infty(\Omega)$

$$y = y(x, t)$$
$$y' = \frac{\partial y}{\partial t}$$

Schrödinger

$$\begin{cases} iy'(x, t) + \operatorname{div}(a(x)\nabla y(x, t)) + q(x)y(x, t) = 0, & x \in \Omega, t \in (0, T) \\ y(x, t) = h(x, t), & x \in \partial\Omega, t \in (0, T) \\ y(x, 0) = y_0(x), & x \in \Omega. \end{cases}$$

Wave

$$\begin{cases} y'' - \operatorname{div}(a(x)\nabla y) + q(x)y = 0, & \Omega \times (0, T) \\ y = h, & \partial\Omega \times (0, T) \\ y(0) = y_0, y'(0) = y_1, & \Omega. \end{cases}$$

If  $a = 1$ ,  $y_0 \in H^1(\Omega)$ ,  $y_1 \in L^2(\Omega)$   $h \in L^2(\partial\Omega \times (0, T))$ ,  $\exists!$  solution.

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## Nonlinear inverse problem :

$$\begin{cases} iy' + \operatorname{div}(a(x)\nabla y) + qy = 0, & \Omega \times (0, T) \\ y = h, & \partial\Omega \times (0, T) \\ y(0) = y_0, & \Omega. \end{cases}$$

Is it possible to retrieve the potential  $q = q(x)$ ,  $x \in \Omega$  from measurement of the normal derivative  $\frac{\partial y}{\partial \nu} \Big|_{\partial\Omega \times (0, T)}$  ?

- ▶  $q \mapsto \frac{\partial y(q)}{\partial \nu}$  is non linear
- ▶ Local result ( $q$  known,  $p$  unknown, close to  $q$ )
- ▶ This is a one-measurement inverse problem.

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# Uniqueness and stability

- ▶ Uniqueness :

$$\left( \frac{\partial y(p)}{\partial \nu} \Big|_{\partial\Omega \times (0,T)} = \frac{\partial y(q)}{\partial \nu} \Big|_{\partial\Omega \times (0,T)} \right) \Rightarrow (p = q \text{ sur } \Omega) ?$$

- ▶ Stability : It is possible to estimate  $(q - p)|_{\Omega}$  by

$$\frac{\partial y(q)}{\partial \nu} - \frac{\partial y(p)}{\partial \nu} \Big|_{\partial\Omega \times (0,T)}$$

in suitable norms ?

# Schrödinger transmission equation

$\Omega_1$ , smooth boundary  $\Gamma_1$  st  $\bar{\Omega}_1 \subset \Omega$  and  $\Omega_0 = \Omega \setminus \bar{\Omega}_1$ .

We set  $a(x) = \begin{cases} a_1 & x \in \Omega_1 \\ a_0 & x \in \Omega_0 \end{cases}$  with  $a_j > 0$  for  $j = 0, 1$

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Inverse problem well-posed provided that :

- ▶  $\Omega_1$  is strongly convex  
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- ▶  $a_0 < a_1$

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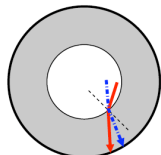
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## Wave transmission equation

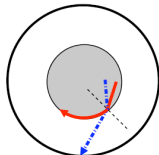
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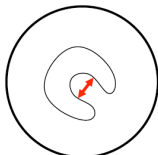
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$a_1 > a_0$  (good)



$a_1 < a_0$  (bad)



$\Omega_1$  non convex (bad)

# References

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- ▶ Wave equation, constant main coefficient

O. Yu. IMANUVILOV, J.-P. PUEL, M. YAMAMOTO 1997, 1999, 2001

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- ▶ I. LASIECKA, R. TRIGGIANI and P. F. YAO 1999

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- ▶ L. CARDOULIS, M. CRISTOFOL and P. GAITAN 2008

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## Method

Takes its roots in the works of Bukhgeim, Klibanov and Malinski.

- ▶ Equation in  $y(q)$
- ▶ Equation in  $z = u'$
- ▶ Equation in  $u = y(q) - y(p)$
- ▶ use of Carleman Estimate

$$\begin{cases} iy' - \operatorname{div}(a\nabla y) + qy = 0, & \Omega \times (0, T) \\ y = h, & \partial\Omega \times (0, T) \\ y(0) = y_0, & \Omega \end{cases} \quad \begin{cases} iu' - \operatorname{div}(a\nabla u) + (f + q)u = fR \\ u = 0 \\ u(0) = 0. \end{cases}$$

$$\begin{cases} u = y(q) - y(p) \\ f = p - q \\ R = y(q) \end{cases} \Rightarrow \text{Linearized version of the inverse problem :}$$

Can we determine  $f$  in  $\Omega$  from  $\frac{\partial u}{\partial \nu} \Big|_{\Omega \times (0, T)}$  ?

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Extension of the equation on  $(-T, T)$

- ▶ Schrödinger : Since  $R(0) \in \mathbb{R}$ , we extend  $v$  and  $R$  on  $\Omega \times (-T, 0)$  by  $v(x, t) = -\bar{v}(x, -t)$  and  $R(x, t) = \bar{R}(x, -t)$ .  
 $z(x, t) = -\bar{z}(x, -t)$ ,  $R(x, t) = \bar{R}(x, -t)$  on  $(-T, 0)$
- ▶ Waves : cut-off function in time and even extension on  $(-T, 0)$

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We take  $q \in L^\infty(\Omega)$  and  $u_0 \in H_0^1(\Omega)$  such that

- ▶  $R = y(q) \in H^1(0, T; L^\infty(\Omega))$
- ▶  $0 < r < |R(x, 0)|$  a.e. on  $\Omega$

and we will use an estimate like

$$C_1(s) \int_{\Omega} |z'(0)|^2 dx \leq C_2(s) \int_{-T}^T \int_{\partial\Omega} \left| \frac{\partial z}{\partial \nu} \right|^2 d\sigma dt + C_3(s) \int_{-T}^T \int_{\Omega} |\text{rhs}|^2 dx$$

with  $\frac{C_3(s)}{C_1(s)} \rightarrow 0$  when  $s \rightarrow \infty$ .

$$\begin{cases} iz' - \text{div}(a\nabla z) + qz = fR', \\ z = 0, \\ z(0) = -ifR(0). \end{cases}$$

# Carleman Estimate for the Schrödinger equation

$Lv = iv' + \Delta v + qv$  we suppose that  $\|q\|_{L^\infty(\Omega)} \leq m$ .

For  $\lambda$  and  $s$  large enough,  $\exists M = M(\Omega, T, m) > 0$  such that :

$$\begin{aligned} & s\lambda \int_{-T}^T \int_{\Omega} |\nabla v|^2 \varphi e^{-2s\varphi} dxdt + s^3 \lambda^4 \int_{-T}^T \int_{\Omega} |v|^2 \varphi^3 e^{-2s\varphi} dxdt \\ & \leq M \int_{-T}^T \int_{\Omega} |Lv|^2 \varphi^3 e^{-2s\varphi} dxdt + Ms\lambda \int_{-T}^T \int_{\partial\Omega} \left| \frac{\partial v}{\partial \nu} \right|^2 d\sigma dt \end{aligned}$$

$\forall v \in L^2(-T, T; H_0^1(\Omega))$  such that  $Lv \in L^2(\Omega \times (-T, T))$  and  $\frac{\partial v}{\partial \nu} \in L^2(\Omega \times (-T, T))$ .

## Waves

$$\phi(x, t) = |x - x_0|^2 - \beta t^2$$
$$\varphi(x, t) = e^{\lambda \phi(x, t)}$$

- no time singularities
- $x_0 \in \mathbb{R}^N \setminus \bar{\Omega}$
- $\beta \in ]0, 1[$

$$Lv = v'' - \Delta v + qv$$

## Schrödinger

$$\varphi(x, t) = \frac{\alpha - e^{\lambda \psi(x)}}{(T-t)(T+t)}$$

- $\lambda > 0$ ,  $\|e^{\lambda \psi}\|_{L^\infty} \leq \alpha$
- $\psi > 0$  regular on  $\Omega$
- $|\nabla \psi| \geq \beta > 0$  in  $\Omega$
- $\exists \varepsilon > 0$  such that  $\forall \xi \in \mathbb{R}^N$ ,  
 $\lambda |\nabla \psi \cdot \xi|^2 + D^2 \psi(\xi, \bar{\xi}) \geq \varepsilon |\xi|^2$
- $\nabla \psi \cdot \nu < 0$  on  $\partial \Omega \setminus \Gamma_0$

$$Lv = iv' + \Delta v + qv$$

$$\text{ex : } \psi(x) = |x - x_0|^2, x_0 \in \mathbb{R}^N \setminus \bar{\Omega}$$

# Proof for the Schrödinger Weight

We set  $w = e^{-s\varphi}v$  and we calculate  $Pw = e^{-s\varphi}L(e^{s\varphi}w)$ .

We get  $Pw = P_1w + P_2w + qw$  with

$$\begin{cases} P_1w = iw' + \Delta w + s^2|\nabla\varphi|^2w \\ P_2w = is\varphi'w + 2s\nabla\varphi \cdot \nabla w + s\Delta\varphi w. \end{cases}$$

$$\begin{aligned} & \int_{-T}^T \int_{\Omega} |Pw - qw|^2 dxdt \\ &= \int_{-T}^T \int_{\Omega} (|P_1w|^2 + |P_2w|^2) dxdt + 2 \operatorname{Re} \int_{-T}^T \int_{\Omega} P_1w \overline{P_2w} dxdt. \end{aligned}$$

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The point is to obtain in the calculus :

- ▶ some “dominant” terms with the right sign

$$s\lambda \int_{-T}^T \int_{\Omega} |\nabla w|^2 dxdt, \quad s^3 \lambda^4 \int_{-T}^T \int_{\Omega} |w|^2 dxdt$$

- ▶ and the “measurement” term

$$s\lambda \int_{-T}^T \int_{\Gamma_0} \left| \frac{\partial w}{\partial \nu} \right|^2 \nabla \psi \cdot \nu d\sigma dt$$

## Discontinuous main coefficient

$$a(x) = \begin{cases} a_1 & x \in \Omega_1 \\ a_0 & x \in \Omega_0 \end{cases} \quad \text{with } a_j > 0 \text{ for } j = 0, 1.$$

Then for each  $f \in L^2(\Omega \times (0, T))$ ,  $y$  satisfies the equation

$$iy' - \operatorname{div}(a\nabla y) + qy = f \quad \text{in } \Omega \times (0, T)$$

iff  $j \in \{0, 1\}$ ,  $y_j = y|_{\Omega_j}$  satisfies

$$iy'_j - a_j \Delta y_j + qy_j = f|_{Q_j} \quad \text{in } \Omega_j \times (0, T)$$

with the *transmission conditions* on the interface  $\Gamma_1 = \partial\Omega_1$

$$\begin{cases} y_0 = y_1 \\ a_0 \frac{\partial y_0}{\partial \nu_0} + a_1 \frac{\partial y_1}{\partial \nu_1} = 0 \end{cases} \quad \text{on } \Gamma_1 \times (0, T).$$

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## Carleman estimate - discontinuous $a$

We use a weight-function  $\varphi(x, t) = \frac{\alpha - e^{\lambda\psi(x)}}{(T-t)(T+t)}$  such that :

- ▶ **Condition 1** :  $\varphi|_{\Omega_j} =$  appropriate weight-function in  $\Omega_j$ .
- ▶ **Condition 2** : We can control the “interface” terms on  $\partial\Omega_1$ .
- ▶ The choice of  $\psi$  is crucial :

$$\begin{cases} |\nabla\psi| \geq \beta > 0 \\ 2D_a^2\psi(\xi, \bar{\xi}) + 2a^2\lambda|\nabla\psi \cdot \xi|^2 - a\nabla a \cdot \nabla\psi|\xi|^2 \geq \epsilon|\xi|^2, \quad \forall \xi \in \mathbb{C}^n \end{cases} \text{ in } \Omega_0 \cup \Omega_1$$

where  $D_a^2\psi \left( a \frac{\partial}{\partial x_i} \left( a \frac{\partial \psi}{\partial x_j} \right) \right)_{1 \leq i, j \leq N}$

$$\begin{cases} \psi_0 = \psi_1 = \text{cst} \\ a_0 \frac{\partial \psi_0}{\partial \nu_0} + a_1 \frac{\partial \psi_1}{\partial \nu_1} = 0 \text{ and } \frac{\partial \psi_0}{\partial \nu_0} + \frac{\partial \psi_1}{\partial \nu_1} < 0 \end{cases} \text{ on } \Gamma_1$$

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## Explicit weight function

We will actually define an  $\varepsilon$ -pair of transmission weight functions  $(\psi^1, \psi^2) \in C^4(\Omega_0 \cup \Omega_1)^2$  such that for each  $j, k \in \{1, 2\}$  with  $j \neq k$ ,

$$\psi^j - \psi^k \geq \delta > 0 \quad \text{in } B_\varepsilon(x_k)$$

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$$\begin{cases} |\nabla \psi^j| \geq \beta > 0 \\ \forall \xi \in \mathbb{C}^n, \\ 2D_a^2 \psi^j(\xi, \bar{\xi}) + 2a^2 \lambda |\nabla \psi^j \cdot \xi|^2 - a \nabla a \cdot \nabla \psi^j |\xi|^2 \geq \epsilon |\xi|^2 \end{cases} \quad \text{in } \Omega_0 \cup \Omega_1 \setminus B_\varepsilon(x_j)$$

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## Intermediate step

$$Z = \left\{ v \in L^2(-T, T; H_0^1(\Omega)) : Lv \in L^2(Q), \frac{\partial v}{\partial \nu} \in L^2(\Sigma), v \text{ satisfies } (Tr) \right\},$$

$$\|w\|_{s,\lambda,\psi} = s^3 \lambda^4 \int_{-T}^T \int_{\Omega} \theta^3 |w|^2 dx dt + s \lambda \int_{-T}^T \int_{\Omega} \theta |\nabla w|^2 dx dt$$

and  $\|\cdot\|_{s,\lambda,\psi,B}$  corresponds to the above terms defined in  $B \subset \Omega$ .

We prove first that

$$\begin{aligned} & \|P_1(w)\|_{L^2}^2 + \|P_2(w)\|_{L^2}^2 + \|w\|_{s,\lambda,\psi}^2 \\ & \leq C \|P(w)\|_{L^2}^2 + C \|w\|_{s,\lambda,\psi,B_\varepsilon(x_j)}^2 + s \lambda C \iint_{\partial\Omega} \left| a \frac{\partial w}{\partial \nu} \right|^2. \end{aligned}$$

## Carleman estimate

Suppose there exists for some  $\varepsilon > 0$  an  $\varepsilon$ -pair of transmission weight functions  $(\psi^1, \psi^2)$  belonging to  $C^4(\Omega_0 \cup \Omega_1)$ .

$\forall q \in L^\infty(\Omega)$  with  $\|q\|_{L^\infty(\Omega)} \leq m$ ,  $\exists C = C(\Omega, T, m) > 0$  such that

$$\begin{aligned} & \sum_{k=1}^2 \left( s\lambda \int_{-T}^T \int_{\Omega} |\nabla v|^2 e^{-2s\varphi^k} dxdt + s^3 \lambda^4 \int_{-T}^T \int_{\Omega} |v|^2 e^{-2s\varphi^k} dxdt \right) \\ & \leq C \sum_{k=1}^2 \left( \int_{-T}^T \int_{\Omega} |Lv|^2 e^{-2s\varphi^k} dxdt + s\lambda \int_0^T \int_{\partial\Omega} \left| a_1 \frac{\partial v}{\partial \nu} \right|^2 d\sigma dt \right) \end{aligned}$$

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# Stability theorem

$$\begin{cases} iy' + \operatorname{div}(a(x)\nabla y) + qy = 0 \\ y = h \\ y(0) = y_0 \end{cases}$$

Let  $\mathcal{U}$  be a bounded subset of  $L^\infty(\Omega)$ ,  $q \in L^\infty(\Omega)$ ,  $T_0 > 0$  and  $r > 0$ .

If  $y_0 \in H_0^1(\Omega)$  and  $y_0$  takes its values in  $\mathbb{R}$  and if

$$\begin{aligned} |y_0(x)| &\geq r, \quad \text{a.e. in } \Omega, \\ y(q) &\in H^1(0, T; L^\infty(\Omega)), \end{aligned}$$

then  $\exists C = C(\Omega_0, \Omega_1, T, a_0, a_1, \|q\|_{L^\infty}, y_0, h, \mathcal{U}, r) > 0$  such that  $\forall p \in \mathcal{U}$ , we have

$$\|q - p\|_{L^2(\Omega)} \leq C \left\| a_1 \frac{\partial y(q)}{\partial \nu} - a_1 \frac{\partial y(p)}{\partial \nu} \right\|_{H^1(0, T; L^2(\Gamma_0))}.$$

$$\begin{cases} iz' - \operatorname{div}(a\nabla z) + qz = fR' \\ z = 0 \\ z(0) = -ifR(0) \end{cases}$$

We set  $w = e^{s\varphi}z$  and define

$$I = \operatorname{Im} \int_{-T}^0 \int_{\Omega} P_1 w \bar{w} \, dx dt.$$

where  $P_1 w = iw' + \Delta w + s^2 |\nabla \varphi|^2 w$ .

► Calculation :  $I = \frac{1}{2} \int_{\Omega} |R(x, 0)|^2 e^{-2s\varphi(0)} |f|^2 \, dx$

► Carleman :

$$I \leq Ms^{-\frac{3}{2}} \left( \int_0^T \int_{\Omega} |fR'|^2 e^{-2s\varphi} \, dx dt + s \int_0^T \int_{\Gamma_0} \left| \frac{\partial v}{\partial \nu} \right|^2 \, d\sigma dt \right)$$

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We finally obtain

$$(1 - Cs^{-\frac{3}{2}}) \int_{\Omega} |f(x)|^2 dx \leq C \int_0^T \int_{\Gamma_0} \left| \frac{\partial v}{\partial \nu} \right|^2 d\sigma dt,$$

and for  $s$  large enough, we prove the stability of the linear inverse problem :

$$\|f\|_{L^2(\Omega)} \leq C \left\| \left\| \frac{\partial u}{\partial \nu} \right\| \right\|_{H^1(0,T;L^2(\Gamma_0))}.$$

## Theorem - Linear inverse problem - Wave equation

Let  $\|q\|_{L^\infty} \leq m$  and  $r > 0$ . We assume that  $a(x) = \begin{cases} a_1 & x \in \Omega_1 \\ a_0 & x \in \Omega_0 \end{cases}$  with

$a_1 > a_0 > 0$  and

$$T > \sqrt{\frac{a_1}{\beta}}$$

$$|R(x, 0)| \geq r \quad \text{a.e. in } \Omega,$$

$$R \in H^1(0, T; L^\infty(\Omega)).$$

Then  $\exists C = C(\Omega_0, \Omega_1, T, \Gamma_0, m, \|R\|) > 0$  such that  $\forall f \in L^2(\Omega)$ ,  
the solution  $u$  of

$$\begin{cases} u'' + \operatorname{div}(a\nabla u) + qu = fR, & \Omega \times (0, T) \\ u = 0, & \partial\Omega \times (0, T) \\ u(0) = 0, u'(0) = 0 & \Omega, \end{cases}$$

satisfies

$$\|f\|_{L^2(\Omega)} \leq C \left\| a_1 \frac{\partial u}{\partial \nu} \right\|_{H^1(0, T; L^2(\partial\Omega))}.$$

## Further problems

- ▶ Recovering the main coefficient.
- ▶ Neumann boundary data and Dirichlet measurement.
- ▶ Recovering two coefficients...

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