

**Control of the Heat Equation for a Discontinuous  
Diffusion Coefficient.**

**Jérôme Le Rousseau and Luc Robbiano**

Université Orléans, UMR CNRS 6628

and

Université de Versailles Saint-Quentin

LMV, UMR CNRS 8100

## THE PROBLEM

Recall : control problem for heat.

We search  $g$  such that the solution of the following problem verifies

$$\left\{ \begin{array}{l} \partial_t u - \Delta u = g \mathbf{1}_\omega \\ u|_{t=0} = u_0 \\ u|_{\partial\Omega} = 0 \end{array} \right.$$

$$u|_{t=T} = 0.$$

$\Omega \subset \mathbb{R}^n$  an open set  $\mathcal{C}^\infty$ .

$\omega \subset \Omega$  an open set.

In this case we say that we have null controllability.

## KNOWN RESULTS

- Russel ('73) (wave  $\Rightarrow$  heat)
- Fursikov & Imanuvilov ('95), Lebeau & Robbiano ('95) (Laplacian with regular coefficients)
- Dobova, Osses & Puel ('02) (transmission problem)
- Fernandez-Cara & Zuazua ('02) (1D, BV coefficient)
- Alessandrini & Escauriaza ('06) (1D,  $L^\infty$  coefficients)
- Benabdallah, Dermenjian & Le Rousseau ('07) (1D, transmission)
- Le Rousseau ('07) ( 1D, BV coefficient, stratified problem multi-D)
- Le Rousseau & Robbiano ('08) ( transmission problem)
- Non linear problem (heat, Navier-Stokes,...)

## NOTATIONS & RESULT

$\Omega$  smooth bounded domain,  $\Omega_1 \subset\subset \Omega$  smooth domain.

$\Omega_2 = \Omega \setminus \bar{\Omega}_1$ ,  $\bar{\Omega}_1 \cap \bar{\Omega}_2 = S$  smooth hypersurface.

$c(x) > 0$  smooth function in  $\Omega_1$  and  $\Omega_2$ , denote  $c_j = c|_{\Omega_j}$ .

$$\Delta_c = \sum_{j=1}^n \partial_{x_j} c \partial_{x_j}.$$

**Theorem 1** *The null controllability is true for heat equation  $\partial_t - \Delta_c$  with Dirichlet boundary condition.*

**Remark 1** *Doubova, Osses and Puel, non linear result but assume*

$$\omega \subset \Omega_1, c_1 < c_2 \text{ on } S$$

$$\omega \subset \Omega_2, c_2 < c_1 \text{ on } S.$$

## CARLEMAN ESTIMATE IN NEIGHBORHOOD OF JUMP

Let  $\psi(x)$ ,  $-\partial_{n_1}\psi|_{\Gamma_1} \gg \partial_{n_2}\psi|_{\Gamma_2}$

$$\Gamma_1 = \Gamma_2 = \partial\Omega_1 \cap \partial\Omega_2 \quad \phi|_{\Gamma_j} = \lim_{x \in \Omega_j, x \rightarrow \Gamma_j} \phi(x)$$

$$\varphi(x) = e^{\lambda\psi(x)} - e^{\lambda K}, \quad K > \sup \psi.$$

**Theorem 2**  $\exists h_0 > 0, C > 0, \forall u$  supported in  $D$ , a neighborhood of  $\Gamma_1$  in  $\Omega$ ,  $u$  satisfying the transmission condition,  $\forall h \in (0, h_0]$

$$h^{-3} \left\| \left( \frac{1}{t(T-t)} \right)^{3/2} e^{\frac{\varphi}{ht(T-t)}} u \right\|^2 + \text{etc.} \leq C \left\| e^{\frac{\varphi}{ht(T-t)}} (\partial_t - \Delta_c) u \right\|^2$$

## SOME REDUCTIONS AND GENERALITIES

-Writing transmission problem like a coupled system at boundary  
(following Bellassoued)

-Conjugate  $v = e^{\varphi/h}u$ , Denote  $P_\varphi = e^{\varphi/h}Pe^{-\varphi/h}$ .

-Transformation  $x_n \rightarrow -x_n$  when  $x_n < 0$ .

-Notations  $v_\ell$  ( $x_n < 0$ ),  $v_r$  ( $x_n > 0$ ), and  $\varphi_\ell, \varphi_r$ .

$$\begin{aligned}
 P_{\varphi_\ell}v_\ell &= f_\ell && \text{in } x_n > 0 \\
 P_{\varphi_r}v_r &= f_r && \text{in } x_n > 0 \\
 v_\ell - v_r &= k_1 && \text{on } x_n = 0 \\
 c_\ell(D_{x_n} + i\frac{1}{t(T-t)}\varphi'_{\ell x_n})v_\ell \\
 + c_r(D_{x_n} + i\frac{1}{t(T-t)}\varphi'_{rx_n})v_r &= k_2 && \text{on } x_n = 0
 \end{aligned} \tag{1}$$

## CARLEMAN APPROACH

-Usual method,  $P_{\varphi_\ell} = A_\ell + iB_\ell$ ,  $P_{\varphi_r} = A_r + iB_r$  ( $A$ ,  $B$  self-adjoint).

$$\|P_{\varphi_\ell} v_\ell\|^2 = \|A_\ell v_\ell\|^2 + \|B_\ell v_\ell\|^2 + 2 \operatorname{Re}(A_\ell v_\ell, iB_\ell v_\ell)$$

-Computation  $\operatorname{Re}(A_\ell v, iB_\ell v)$  gives usual terms and boundary terms.

$$\|P_{\varphi_\ell} v_\ell\|^2 + \|P_{\varphi_r} v_r\|^2 = U(v) + \mathcal{B}(v)$$

-where  $U(v)$  are the usual wellknown terms, “easy”.

- $\mathcal{B}(v)$  are terms on boundary like  $(L_0 D_{x_n} v_r, D_{x_n} v_r)_0$ ,  $(L_1 v_r, D_{x_n} v_r)_0$ ,  $(L_2 v_r, v_r)_0$ ,  $L_j$  are differential tangential operators of order  $j$ . ( $(\cdot, \cdot)_0$  means inner product on boundary).

-We can use boundary conditions to analyse  $\mathcal{B}(v)$  but in general  $\mathcal{B}(v) \not\geq 0$ . In case Doubova-Osses-Puel  $\mathcal{B}(v) \geq 0$

## MODEL PROBLEM

“Formal” solution  $\Delta u = 0$  in  $x_n > 0$  with boundary conditions.

Fourier with respect  $x' = (x_1, \dots, x_{n-1})$ , we have

$$(\partial_{x_n}^2 - |\xi'|^2)\hat{u}(\xi', x_n) = 0.$$

The solution  $\hat{u}$  is given by

$$\hat{u}(\xi', x_n) = A(\xi')e^{-x_n|\xi'|} + B(\xi')e^{x_n|\xi'|}.$$

but  $e^{x_n|\xi'|} \notin \mathcal{S}'$  if  $x_n > 0 \rightarrow$  “false solution” then

$$\hat{u}(\xi', x_n) = A(\xi')e^{-x_n|\xi'|}.$$

In particular  $\hat{u}(\xi', 0) = A(\xi')$ ,  $\partial_{x_n}\hat{u}(\xi', 0) = -|\xi'|A(\xi')$  then

$$|\xi'|\hat{u}(\xi', 0) + \partial_{x_n}\hat{u}(\xi', 0) = 0$$

It follows : **one** trace determine the **two** traces.

Dirichlet-to-Neumann, Neumann-to-Dirichlet etc.

## ELLIPTIC BOUNDARY PROBLEM

$$Pv = f \text{ in } x_n > 0$$

$$\text{Denote } \underline{v} = \begin{cases} v & \text{if } x_n > 0 \\ 0 & \text{if } x_n < 0 \end{cases}$$

$$P\underline{v} = \underline{f} + \gamma_0 \delta'_{x_n=0} + \gamma_1 \delta_{x_n=0}$$

where  $\gamma_j$  depend of  $v|_{x_n=0}$  and  $D_{x_n} v|_{x_n=0}$ .

Choose  $Q$  a parametrix of  $P \longrightarrow QP = Id + R$  where  $R$  smoothing.

$\underline{v} = Q\underline{f} + C_0 \gamma_0 + C_1 \gamma_1 + R\underline{v}$  where  $C_j$  are operators on traces.

On  $x_n = 0^+$

$$v|_{x_n=0} = Q\underline{f}|_{x_n=0} + C'_0 \gamma_0 + C'_1 \gamma_1 + R\underline{v}|_{x_n=0}$$

Does this give relation between  $v|_{x_n=0}$  and  $D_{x_n} v|_{x_n=0}$ ?

## CALDERON OPERATORS I

$C_0, C_1$  have the following form (principal symbol)

$$\int e^{ix'\xi'} \hat{\gamma}(\xi') \int e^{ix_n\xi_n} \frac{\xi_n^\mu}{p(x, \xi)} d\xi_n d\xi'$$

**KEY POINT**  $\longrightarrow$  roots in  $\xi_n$  of  $p(x, \xi', \xi_n)$ .

**WHY?**  $\longrightarrow$  residues formula.

---

$p(x, \xi', \xi_n)$  is the principal symbol of  $P$ .

$\xi_n^\mu$  comes from Fourier transform of  $\delta_{x_n=0}^{(\mu)}$ .

## CALDERON OPERATORS II

Let  $p(x, \xi) = (\xi_n - r_1(x, \xi'))(\xi_n - r_2(x, \xi')) \rightarrow 3$  cases

- If  $\text{Im } r_1(x, \xi') < 0$  and  $\text{Im } r_2(x, \xi') < 0$  then

$$\int e^{ix_n \xi_n} \frac{\xi_n^\mu}{p(x, \xi)} d\xi_n = 0$$

$f$  determine the traces of  $u$  (analogous computation for  $\partial_n u$ ).

- If  $\text{Im } r_1(x, \xi') > 0$  and  $\text{Im } r_2(x, \xi') < 0$  then

$$\int e^{ix_n \xi_n} \frac{\xi_n^\mu}{p(x, \xi)} d\xi_n = 2i\pi e^{ix_n r_1(x, \xi')} \frac{r_1^\mu(x, \xi')}{r_1(x, \xi') - r_2(x, \xi')}$$

We find the model case and we have a relation between the two traces

- If  $\text{Im } r_1(x, \xi') > 0$  and  $\text{Im } r_2(x, \xi') > 0$  no relations between the two traces

## BOUNDARY CONDITIONS AND ROOTS LOCALISATION

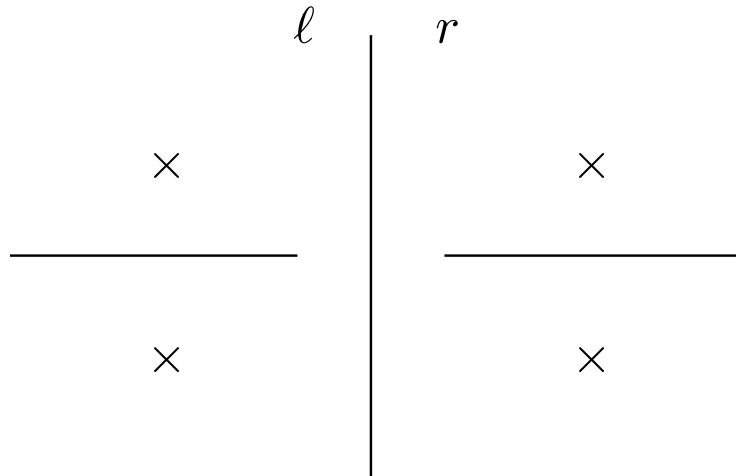
×  
 \_\_\_\_\_ One relation between the two traces  $u|_{x_n=0}$  and  $D_{x_n}u|_{x_n=0}$ .  
 ×

\_\_\_\_\_ The two traces  $u|_{x_n=0}$  and  $D_{x_n}u|_{x_n=0}$  are determined.  
 × ×

× \_\_\_\_\_ Carleman method  $\longrightarrow \simeq$  One relation between the two traces.  
 ×

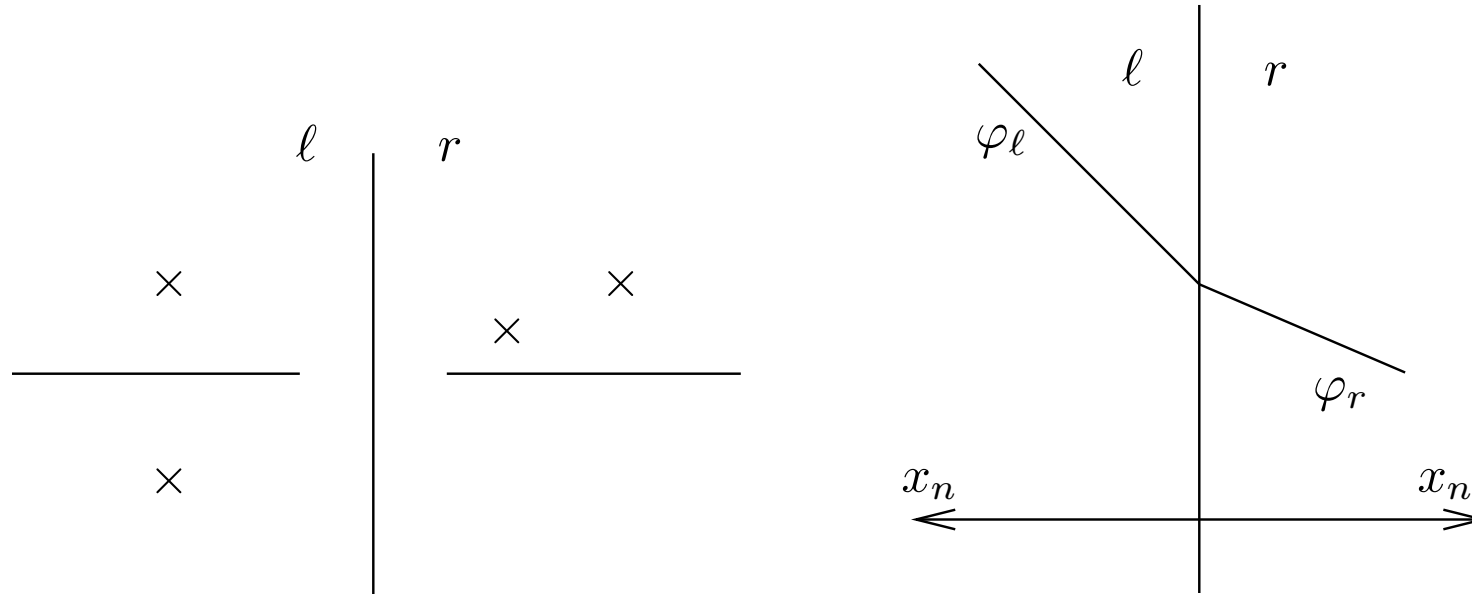
×  
 × ×  
 × \_\_\_\_\_ No relation between two traces  $u|_{x_n=0}$  and  $D_{x_n}u|_{x_n=0}$

## RETURN TO COUPLED BOUNDARY PROBLEM



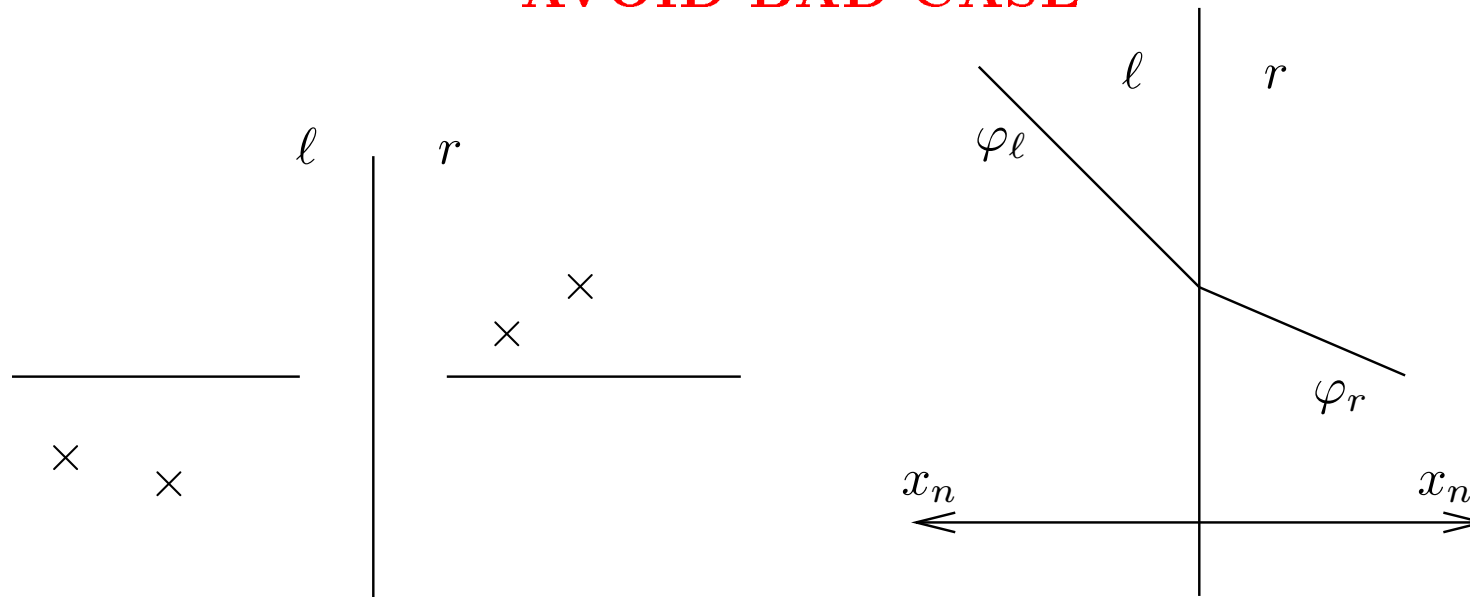
- This case exists if  $|\xi'|$  large.
  - One relation between the traces at left.
  - One relation between the traces at right.
  - Two transmission conditions.
- 4 equations for 4 unknown and ... OK (algebraic checking)

## A BAD CASE



- One relation between the traces at left.
  - No condition between the traces at right.
  - Two transmission conditions.
- 3 equations for 4 unknown !!!! Impossible to solve.

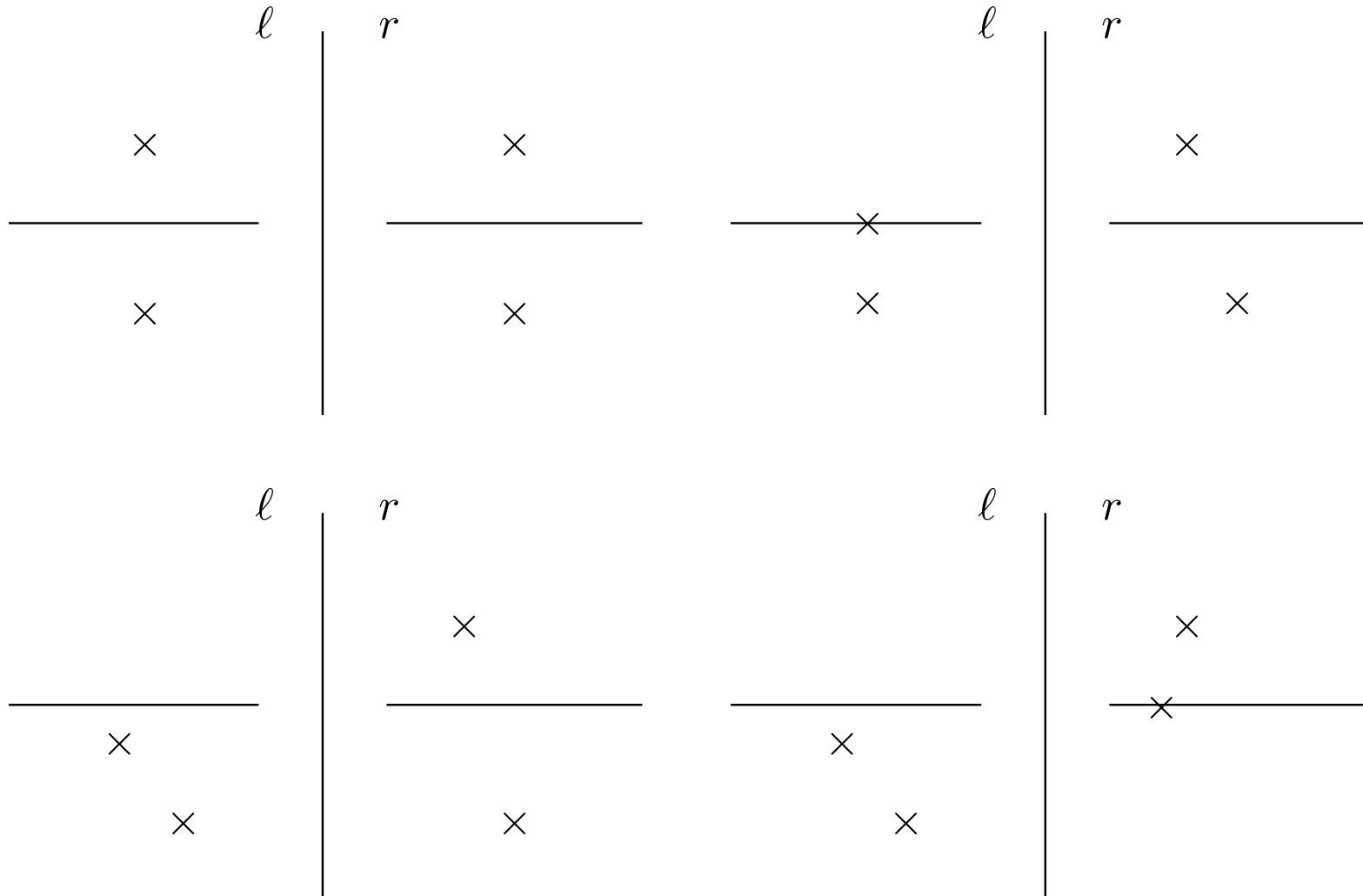
## AVOID BAD CASE



Choose parameters such that the two roots at the left are in  $\text{Im } z < 0$  when the two roots at the right are in  $\text{Im } z > 0$ .

- The traces at left are known  $\simeq$  two equations
- No condition between the traces at right.
- Two transmission conditions  
→ 4 equations for 4 unknown and ... OK.

## DIFFERENT CASES, MOVIE



## COMPUTATIONS WITH FURSIKOV-IMANUVILOV WEIGHTS

Computation  $e^{\frac{\varphi}{ht(T-t)}} P(e^{\frac{-\varphi}{ht(T-t)}} u)$  où  $P = h^2(\partial_t - \Delta)$ .

Symbol rule  $\xi \leftrightarrow \frac{h}{i}\partial_x$  et  $\tau \leftrightarrow \frac{h^2}{i}\partial_t$ .

Symbol  $= i\tau + |\xi|^2 + 2i\frac{\xi\varphi'}{t(T-t)} - \frac{|\varphi'|^2}{t^2(T-t)^2} - h\frac{(T-2t)\varphi}{t^2(T-t)^2}$ .

Recall : we want distinguish elliptic zone and “Carleman” zone.

→ Symbolic calculus in  $\tau$  and  $\xi$ .

Each term must be of the same “strength”  $\tau$ ,  $|\xi|^2$  et  $\frac{1}{t^2(T-t)^2}$ .

Weight  $\left( |\tau| + |\xi|^2 + \frac{1}{t^2(T-t)^2} \right)^{1/2}$

“Bad symbolic calculus ” in Hörmander sense : singularities in  $t = 0$   
et  $t = T$ .

## GOOD SYMBOLIC CALCULUS

Transformation  $s(t) = \tan\left(\frac{\pi t}{T} - \frac{\pi}{2}\right)$

goal  $t \rightarrow T^-$  then  $s \rightarrow +\infty$  and  $t \rightarrow 0^+$  then  $s \rightarrow -\infty$ .

$\partial_t \rightarrow a(s)\partial_s$  where  $a(s) \sim \langle s \rangle^2$

$\eta(t) = \frac{1}{t(T-t)}$  in variable “ $s$ ” we have  $\eta(s) \sim \langle s \rangle$ .

Symbol  $= ia(s)\tau + |\xi|^2 + 2i\xi\varphi'\eta(s) - |\varphi'|^2\eta^2(s) - h\varphi\eta^2(s)$ .

Roughly speaking  $i\langle s \rangle^2\tau + |\xi|^2 + 2i\xi\langle s \rangle - \langle s \rangle^2$

Adapted weight  $\Lambda^4 = \langle s \rangle^4 + \langle s \rangle^4\tau^2 + |\xi|^4$ .

Symbol of order  $\Lambda^2$ .

Calculus in  $S(\Lambda^2, g)$  where  $g = \frac{ds^2}{\langle s \rangle^2} + dx^2 + \frac{\langle s \rangle^4 d\tau^2}{\Lambda^4} + \frac{d\xi^2}{\Lambda^2}$ .

Good analysis for elliptic points.

## CARLEMAN AND SYMBOLIC CALCULUS

Fact : previous calculus are not adapted for Carleman method.

Real part =  $|\xi|^2 - |\varphi'|^2 \eta^2(s) - h\varphi \eta^2(s)$ .

Roughly speaking  $|\xi|^2 - \langle s \rangle^2$ .

Imaginary part =  $a(s)\tau + 2\xi\varphi'\eta(s)$ . Roughly speaking  $\langle s \rangle^2\tau + \xi\langle s \rangle$

We must balance the two terms of real part and imaginary part.

Goal search a weight giving the same “strength” to  $|\xi|^2$  and  $\langle s \rangle^2$

and the same “strength” to  $\langle s \rangle^2\tau$  and  $\langle s \rangle\xi$

Adapted weight :  $M^2 = \langle s \rangle^2 + \langle s \rangle^2\tau^2 + |\xi|^2$ .

Good calculus  $S(M^2, \tilde{g})$  where  $\tilde{g} = \frac{ds^2}{\langle s \rangle^2} + dx^2 + \frac{\langle s \rangle^2 d\tau^2}{M^2} + \frac{d\xi^2}{M^2}$ .

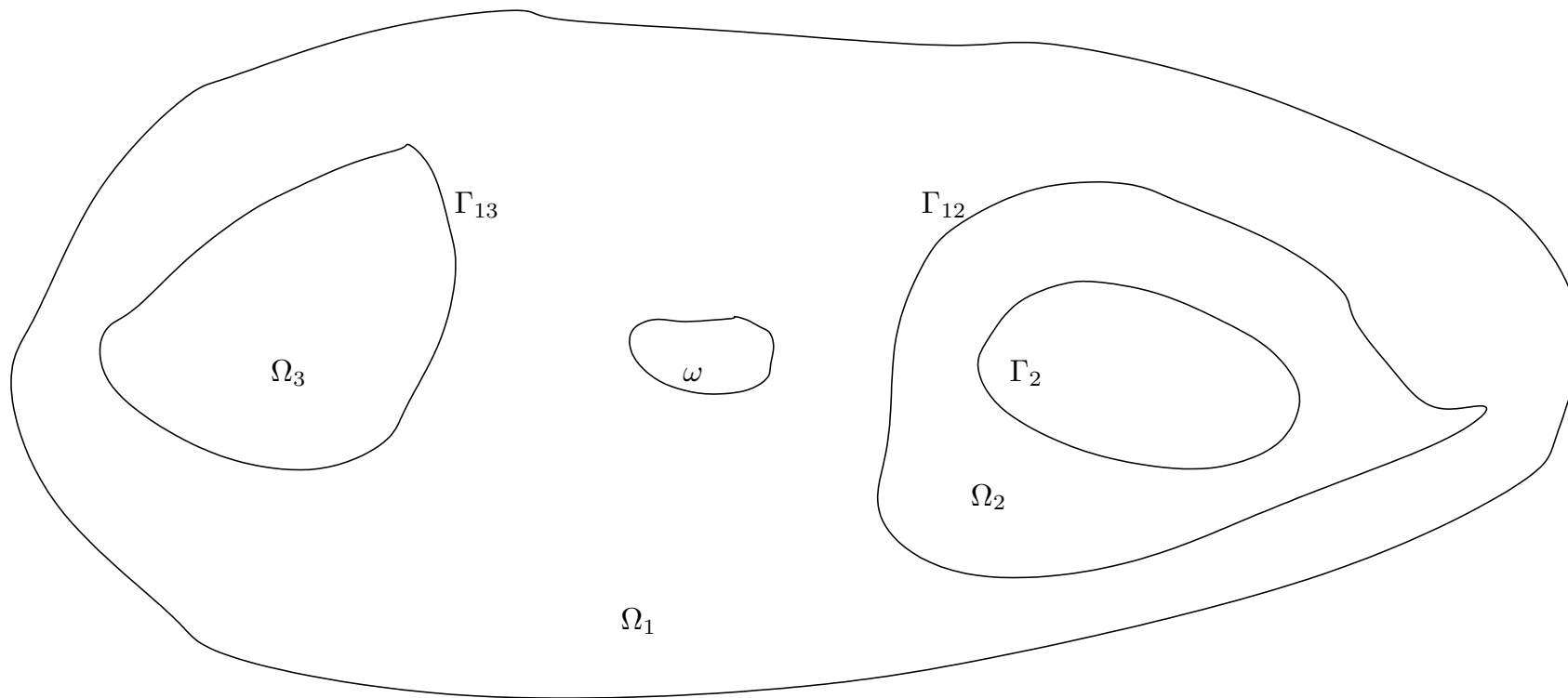
Real part of order  $M^2$ , imaginary part of order  $\langle s \rangle M$ .

The Poisson’s bracket (symbol of the commutator) : order  $\langle s \rangle M^2$ .

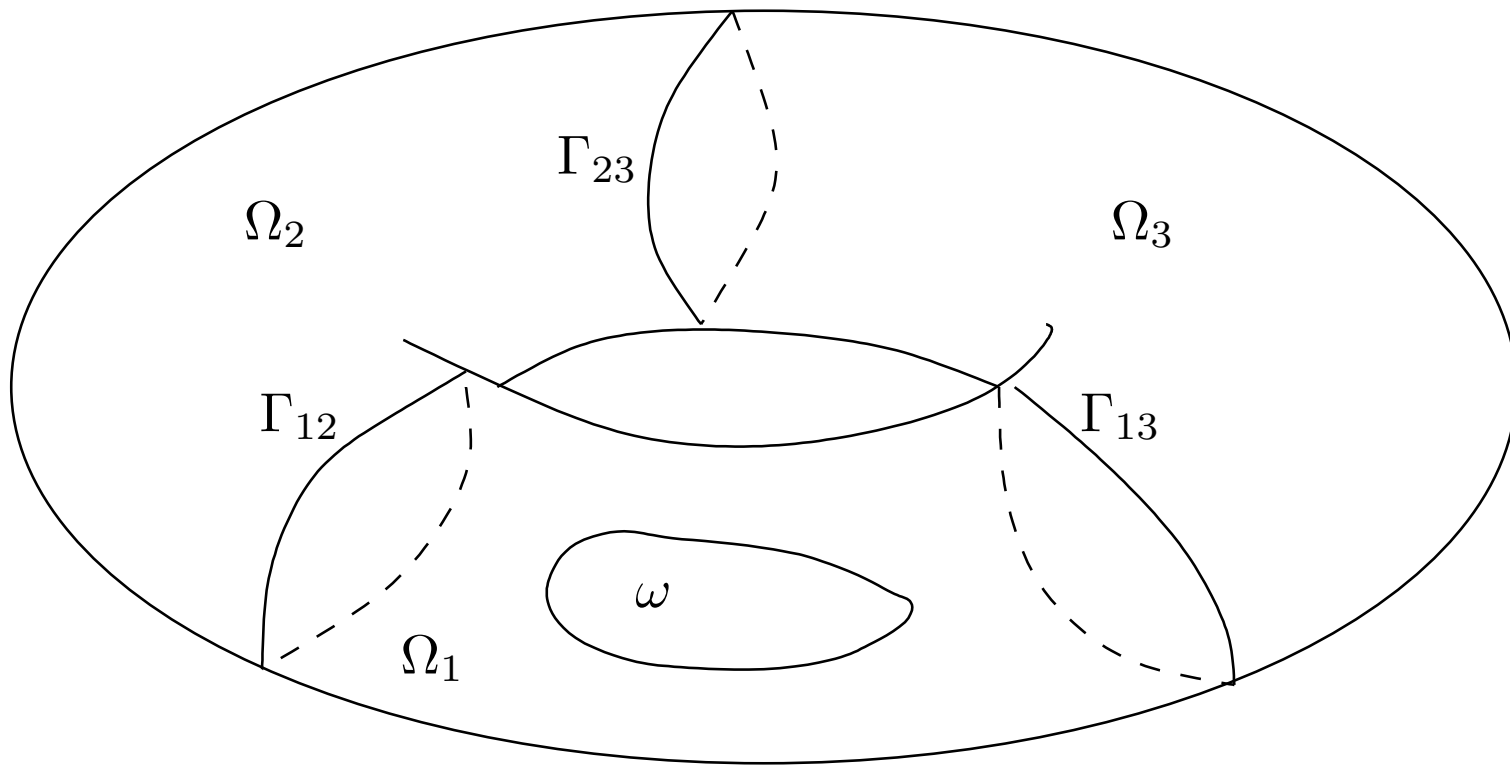
## **SOME PROBLEMS TO FINISH THE CALCULUS**

- Managed to kind of estimates, elliptic zone and “Carleman” zone.
- Glued estimates  $\rightarrow$  need several weight.

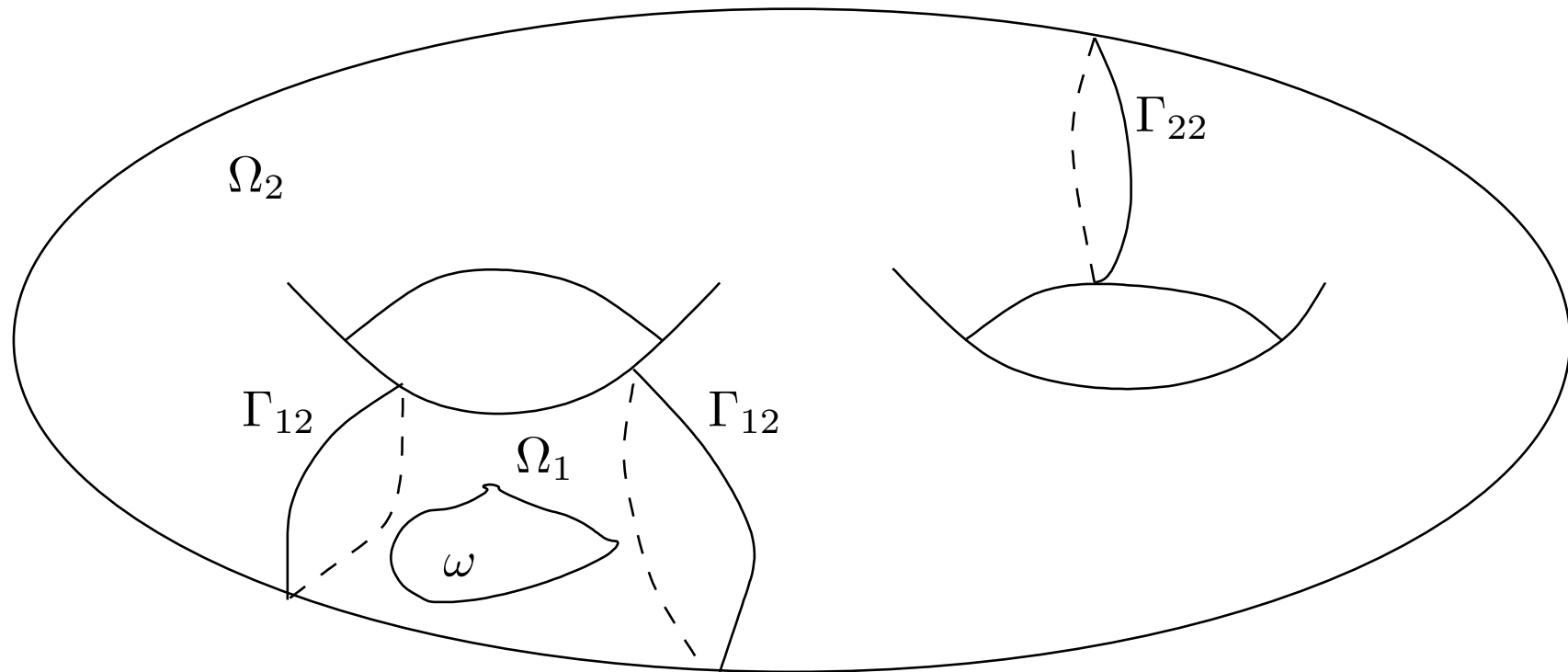
**EXAMPLE : AN OPEN**



# TORUS WITH ONE HOLE



# TORUS WITH 2 HOLES



# OPEN PROBLEM

